



Spring and Damper in Modelling

#8



Serdar ARITAN

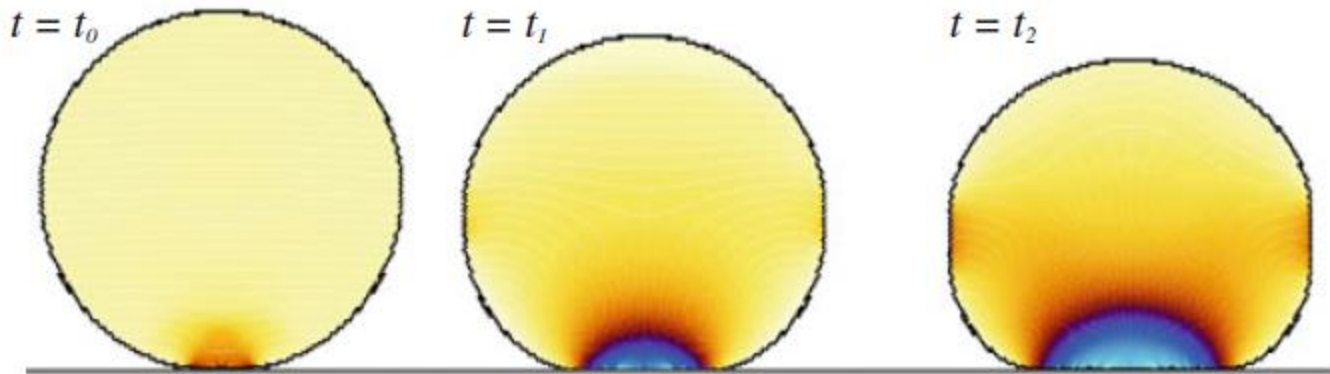
Biomechanics Research Group,
Faculty of Sports Sciences, and
Department of Computer Graphics
Hacettepe University, Ankara, Turkey



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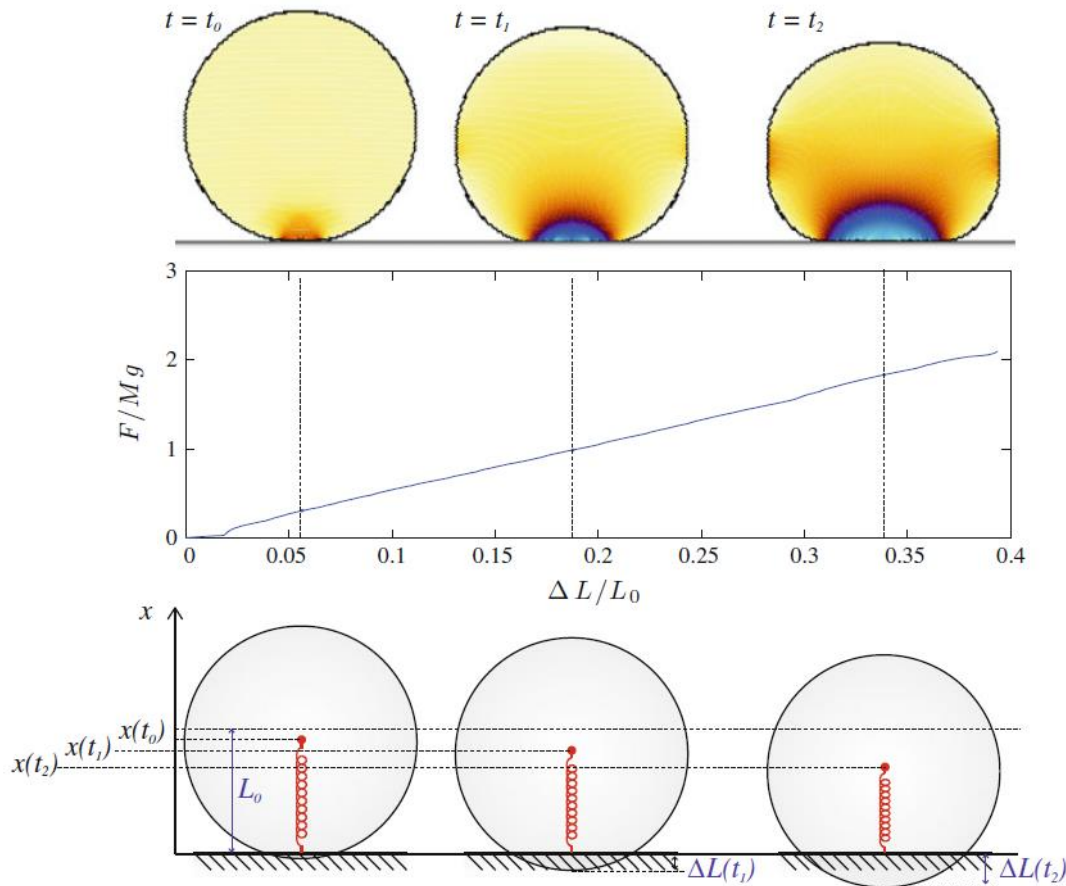
Contact Forces

The spring force was introduced as an approximative model for the force due to deformation. It is based on experimental evidence: We find the law by measuring the force as a function of the deformation. And the law is surprisingly versatile: We can use it as a model for almost any contact force between solid objects.





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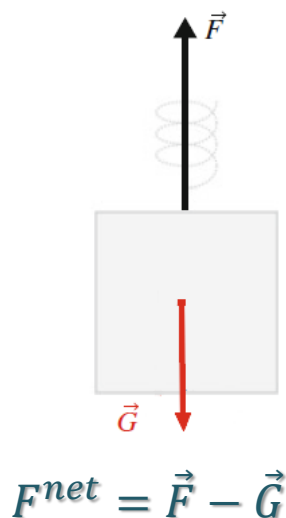
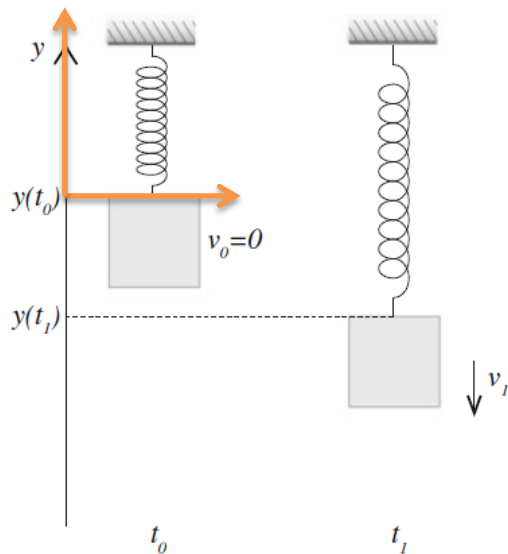
we use the **spring model** as an approximative model for the contact force. Generally, we do not know what the spring constant will be for such a model. We need either to measure the spring constant or to find the spring constant from a theoretical consideration based on for example elasticity theory.



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Motion of a Hanging Block

A block of mass $m = 1$ kg is hanging from a spring with spring constant $k = 100$ N/m. The other end of the spring is attached to the ceiling. We apply the structured problem-solving approach to find the motion of the block after it is released.



$$\vec{F} = \pm k\Delta L$$

For simplicity, we place the coordinate system so that the spring is in equilibrium when $y = 0$ m

$$\vec{F}(y) = -ky\mathbf{j}$$



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$$F^{net} = \vec{F} - \vec{G}$$

$$ma = -ky - mg$$

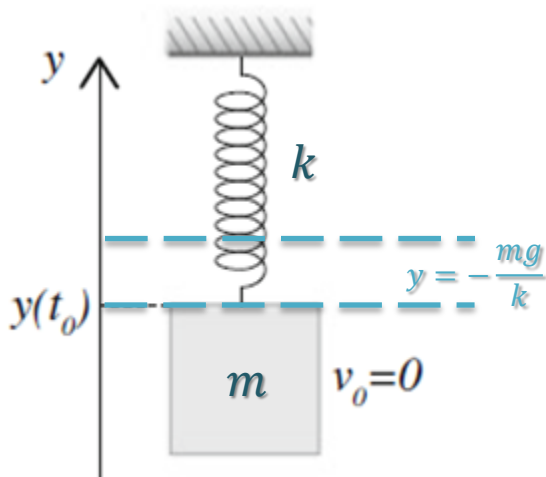
Equilibrium model: First, let us consider the equilibrium situation—where the block does not move when released. $a = 0$

$$ma = -ky - mg = 0$$

$$ma + ky + mg = 0$$

$$ky = -m(a + g)$$

$$y = -\frac{mg}{k}$$





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Numerical solution: We can find the motion of the block using an Euler scheme. Generally, we advice you to use a **fourth-order Runge-Kutta** method for oscillator problems, but we use Euler here to make the programming transparent.

```
import numpy as np
import matplotlib.pyplot as plt
```

```
# Initialize
```

```
m = 1.0      # kg
```

```
k = 100.0    # N/m
```

```
v0 = 1.0     # m/s
```

```
time = 2.0   # s
```

```
g = 9.8      # m/s^2
```

```
# Numerical setup
```

```
dt = 0.0001 # s
```



time step: please try 60 Hz (1/60)

```
n = int(round(time/dt))
```

```
t = np.zeros(n, float)
```

```
y = np.zeros(n, float)
```

```
v = np.zeros(n, float)
```



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```
# Initial values
```

```
y[0] = 0.0
```

```
v[0] = v0
```

```
# Simulation loop
```

```
for i in range(n-1):
```

```
    F = -k*y[i] - m*g
```

```
    a = F/m
```

```
    v[i+1] = v[i] + a*dt
```

```
    y[i+1] = y[i] + v[i+1]*dt
```

```
    t[i+1] = t[i] + dt
```

```
fig, (ax1, ax2) = plt.subplots(2, 1)
```

```
ax1.plot(t,y,'-b')
```

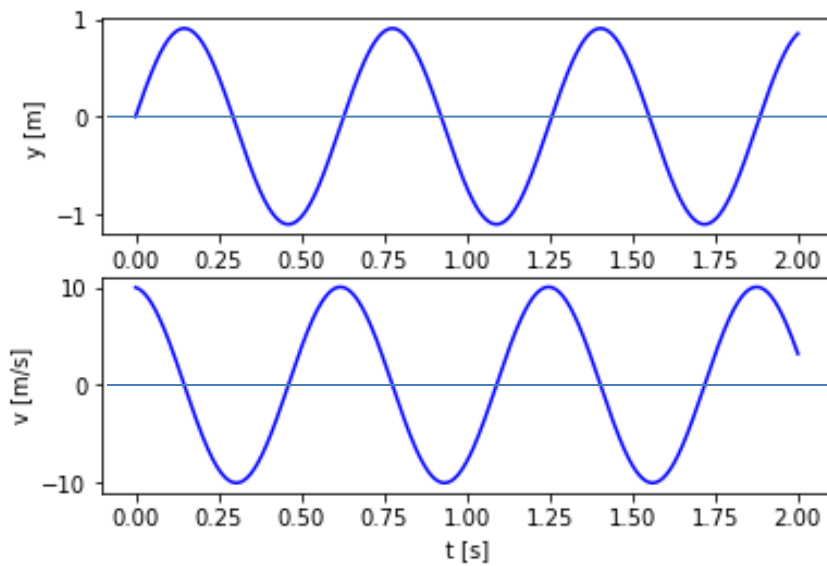
```
ax2.plot(t,v,'-b')
```

```
ax1.set_ylabel('y [m]')
```

```
ax2.set_ylabel('v [m/s]')
```

```
ax2.set_xlabel('t [s]')
```

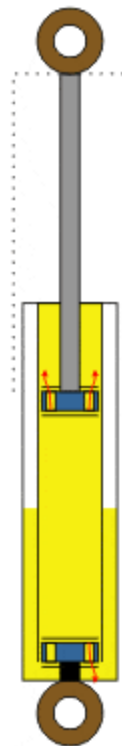
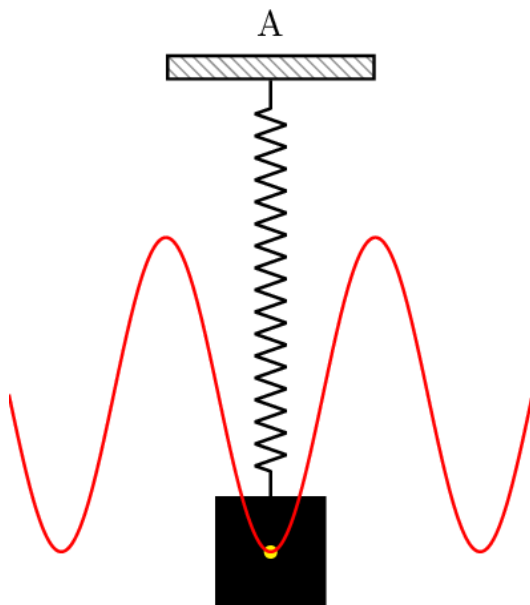
```
plt.show()
```





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Spring and Damper





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Spring and Damper

Spring: an ideal elastic element

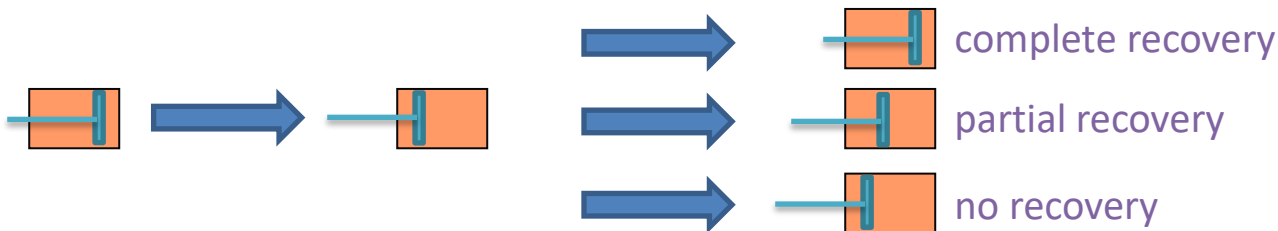
Will immediately change and return to its original shape upon loading and unloading.



Damper: a viscous fluid

Will change its original shape upon loading, depending on time (and temperature).

May slowly return to or may not return to its original shape upon unloading.
partial or complete recovery

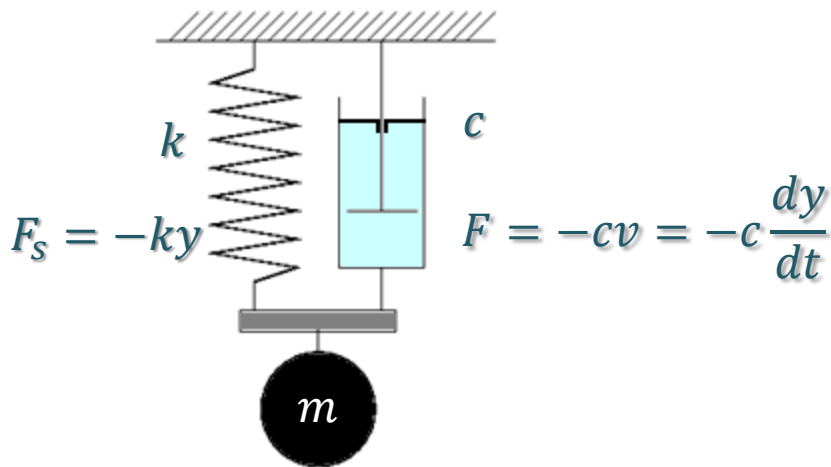




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Mass-Spring-Damper (MSD)

An ideal mass-spring-damper system with mass m , spring constant k , and viscous damper of damping coefficient c is subject to an oscillatory force.



$$F^{net} = -ky + \left(-c \frac{dy}{dt} \right)$$

$$ma = -ky - c \frac{dy}{dt}$$

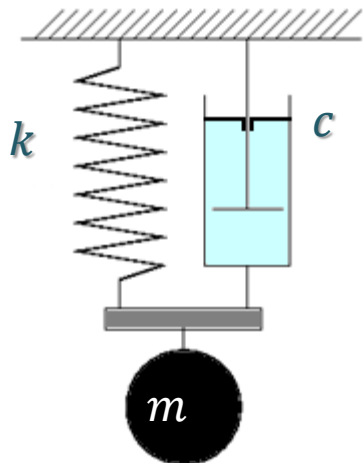
$$a = -\frac{k}{m}y - \frac{c}{m} \frac{dy}{dt}$$

$$\frac{d^2y}{dt^2} + \frac{c}{m} \frac{dy}{dt} + \frac{k}{m}y = 0$$

2. Order Differential Equation



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$$\zeta = \frac{c}{2\sqrt{km}}$$

the damping ratio

$$\omega_0 = \sqrt{\frac{k}{m}}$$

the natural frequency of the system

$$\frac{d^2y}{dt^2} + 2\zeta\omega_0 \frac{dy}{dt} + \omega_0^2 y = 0$$

The first parameter, ω_0 , is called the (undamped) natural frequency of the system. The second parameter, ζ , is called the damping ratio. The natural frequency represents an angular frequency, expressed in radians per second. The damping ratio is a dimensionless quantity.



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By choosing appropriate parameters we can “tune” the behavior of the system

We can classify the system into three distinct categories.

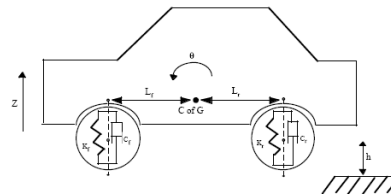
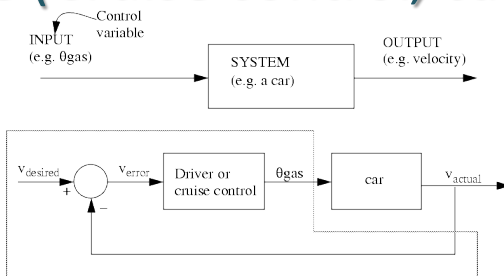
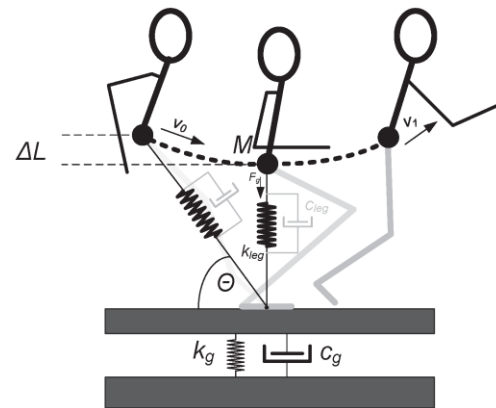
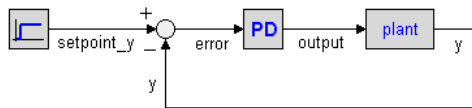
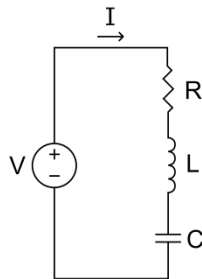
1. Under-damped: $\zeta < 1$. In this case the system oscillates with a frequency equal to $\omega_d = \omega_0 \sqrt{1 - \zeta^2}$
2. Over-damped: $\zeta > 1$. The system slowly returns to equilibrium
3. Critically Damped: $\zeta = 1$. The system returns to equilibrium



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MSD-like Systems are found throughout engineering

- Human Running
- Circuit Design: RLC circuits
- Robotics (PD control)
- Automotive (Cruise control, car suspensions etc.)





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```
import numpy as np
import matplotlib.pyplot as plt
# Initialize
timeFinal= 16.0    # Simulation time in seconds
steps = 10000      # Number of steps
dt = timeFinal/steps    # Step length => 0.0016 s
t = np.linspace(0, timeFinal, steps+1)
# Creates an array with steps+1 values from 0 to timeFinal
# Allocating arrays for velocity and position
v = np.zeros(steps+1)
y = np.zeros(steps+1)
# Setting constants and initial values for vel. and pos.
k = 0.1
m = 0.01
v0 = 0.05
y0 = 0.01
freqNatural = (k/m)**0.5
c = 0.00    #Undamped
v[0] = v0    #Sets the initial velocity
y[0] = y0    #Sets the initial position
```



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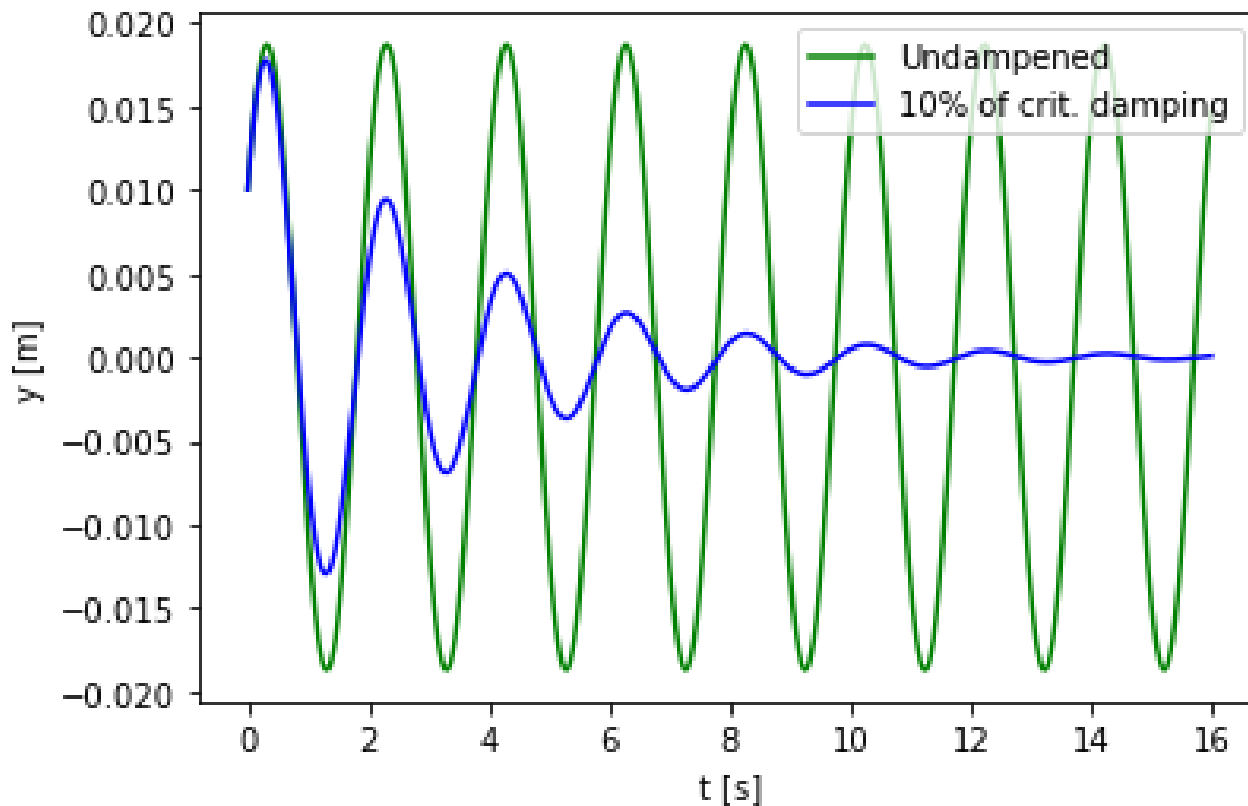
```
# Simulation loop
# Numerical solution using Euler's
#  $v'(t) = -(k/m)*x(t) - (c/m)*v(t)$ 
#  $x'(t) = v(t)$ 
for i in range(0, steps):
    v[i+1] = (-k/m)*dt*y[i] + v[i]*(1-dt*c/m)
    y[i+1] = y[i] + v[i+1]*dt
fig, ax = plt.subplots()
ax.plot(t, y, '-g', label='Undamped')
# Damping set to 10% of critical damping
c = (2*(k*m)**0.5)*0.10
for i in range(0, steps):
    v[i+1] = (-k/m)*dt*y[i] + v[i]*(1-dt*c/m)
    y[i+1] = y[i] + v[i+1]*dt

ax.plot(t, y, 'b-', label = '10% of crit. damping')
ax.xlabel('t [s]')
ax.set_ylabel('y [m]')
ax.legend(loc = 'upper right')
plt.show()
```




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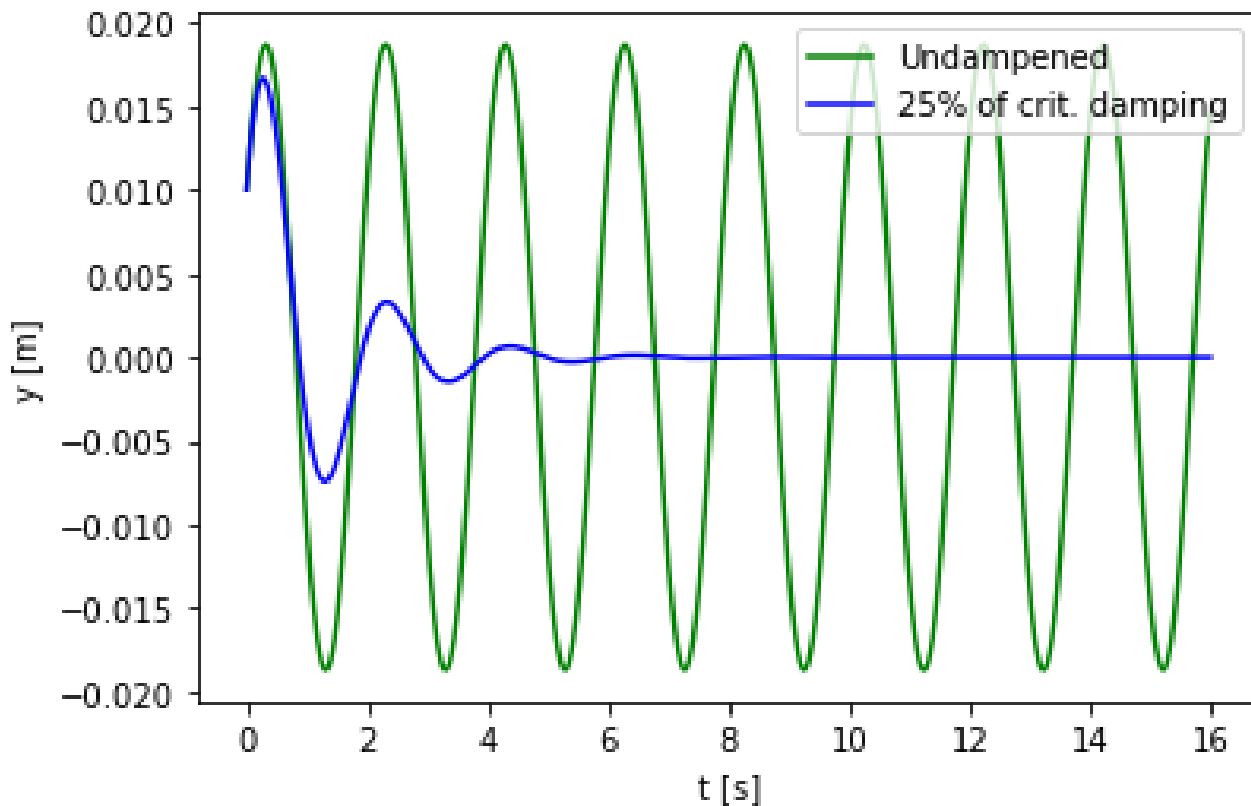
$$c = (2 * (k * m) ** 0.5) * 0.10$$





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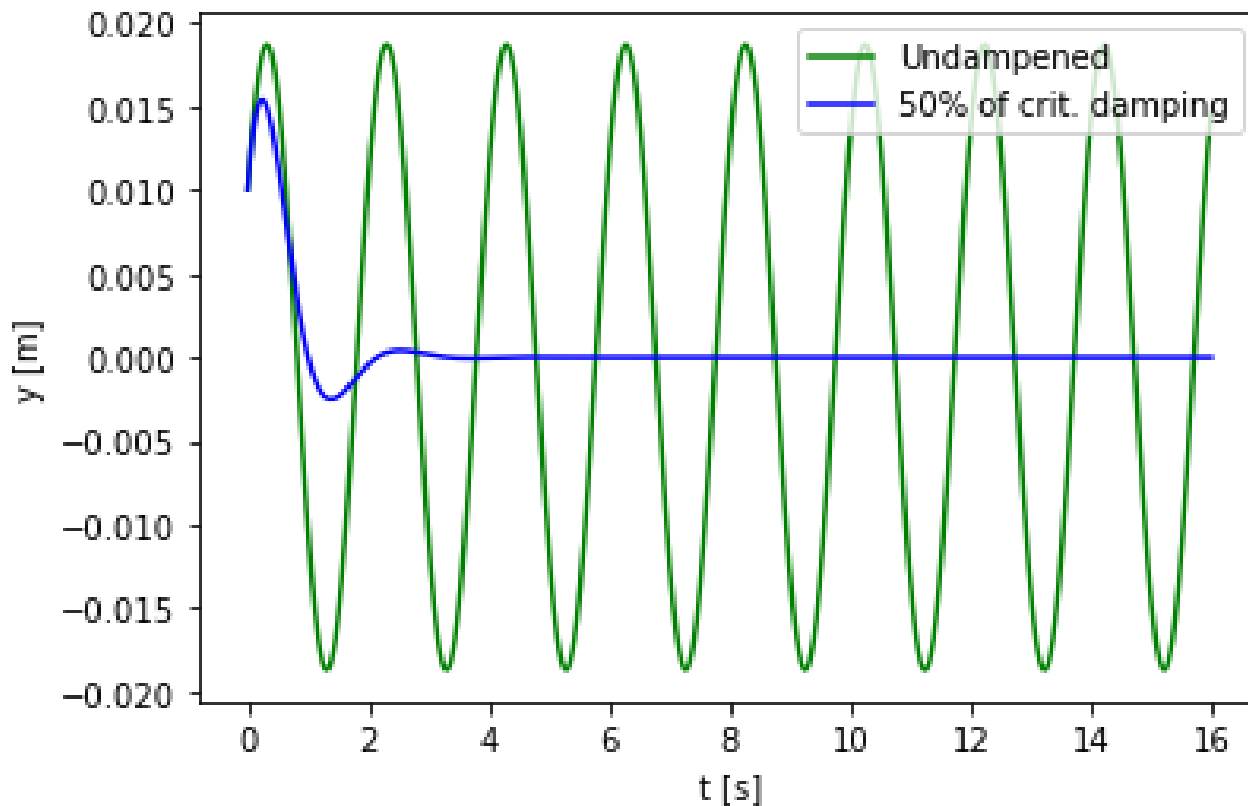
$$c = (2 * (k * m) ** 0.5) * 0.25$$





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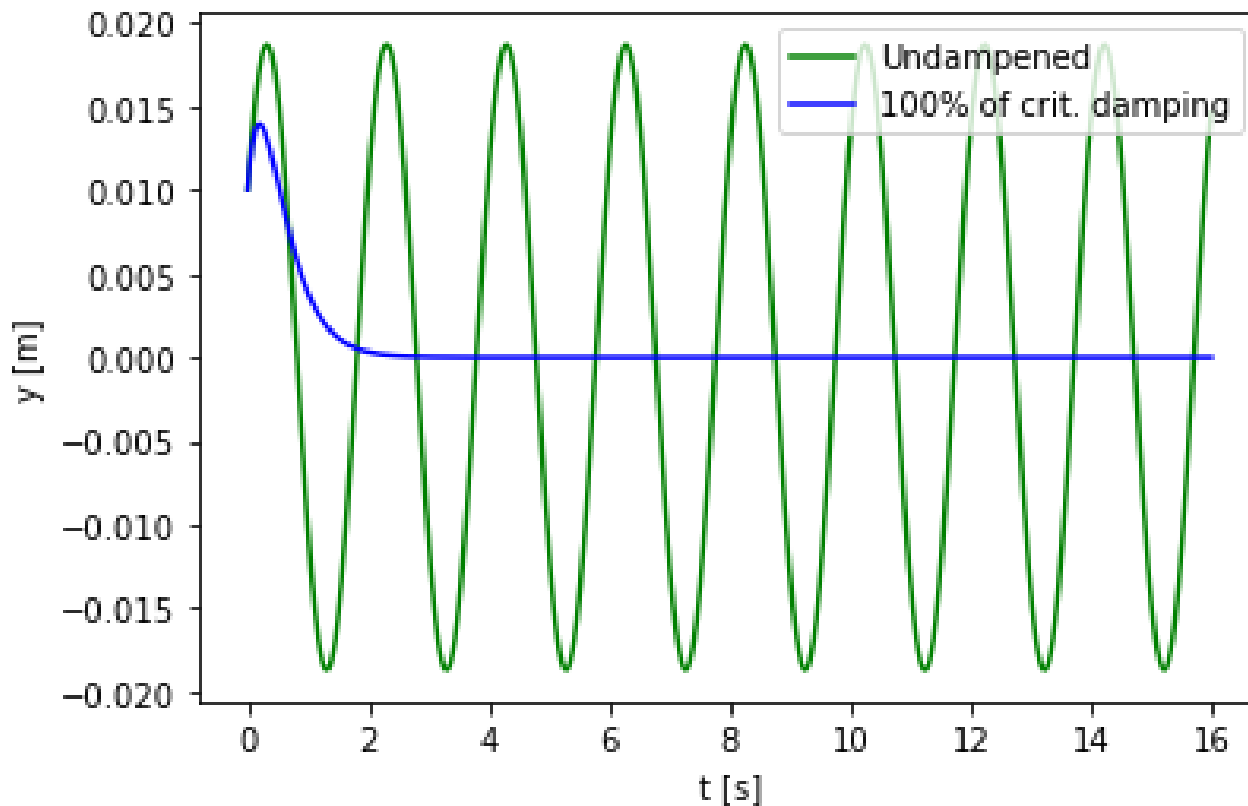
$$c = (2 * (k * m) ** 0.5) * 0.50$$





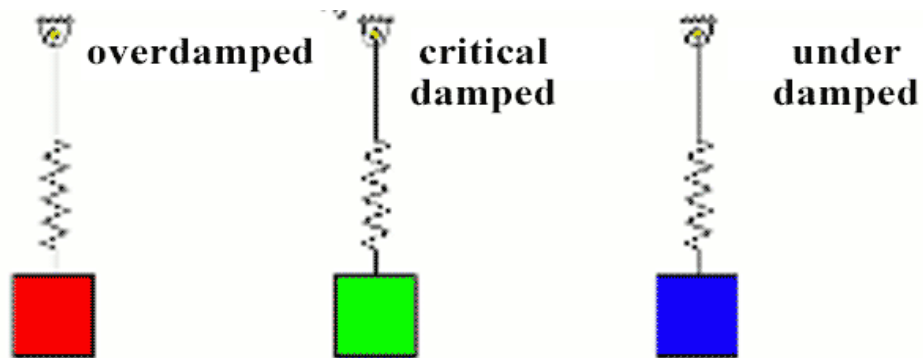
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$$c = (2 * (k * m) ** 0.5) * 1.00$$



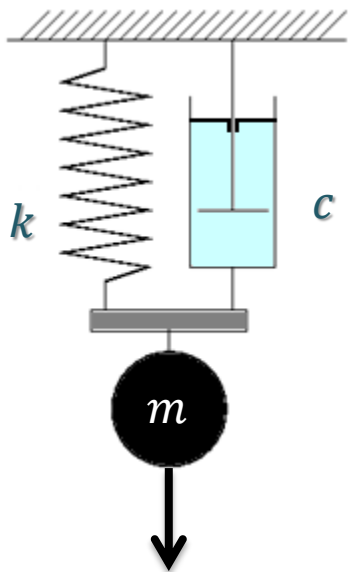


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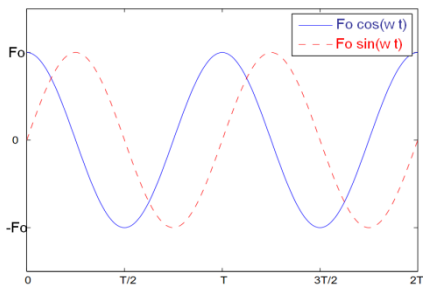




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$$F(t) = F_0 \sin \omega t$$



External Forcing

External Forcing models the behavior of a system which has a time varying force acting on it. An example might be a suspended bridge subjected to wind loading.





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Natural frequency is the frequency at which a system tends to oscillate in the absence of any driving or damping force.

Natural frequency depends on the elasticity and shape of the object.

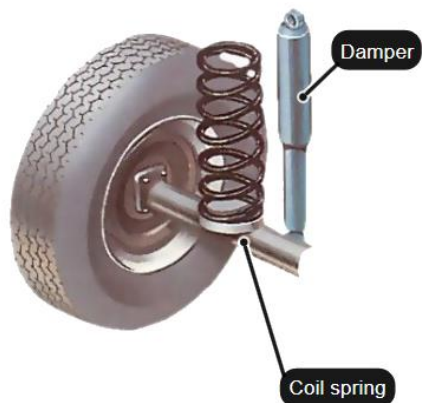
The natural frequency of the smaller bell is higher than that of the big bell, and it rings at a higher pitch.



The natural frequency of a trombone can be modified by changing the length of the air column inside the metal tube.



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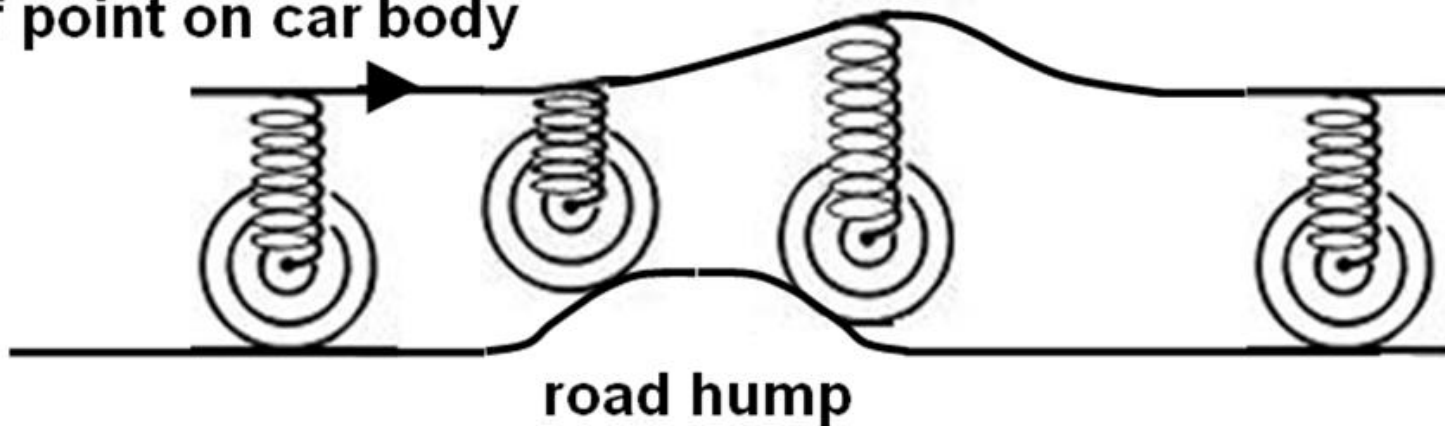


$$F(t) = F_0 \cos \omega t$$

$$F_0 = 0.01$$

$$\omega = \sqrt{k/m}$$

path of point on car body





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```
import numpy as np
import matplotlib.pyplot as plt

# Initialize
timeFinal= 16.0    # Simulation time in seconds
steps = 10000      # Number of steps
dt = timeFinal/steps    # Step length
t = np.linspace(0, timeFinal, steps+1)
# Creates an array with steps+1 values from 0 to timeFinal
v = np.zeros(steps+1)
y = np.zeros(steps+1)
# Setting constants and initial values for vel. and pos.
k = 0.1
m = 0.01
v0 = 0.01
y0 = 2.00
freqNatural = (k/m)**0.5
c = 0.00    #Undamped
F0 = 0.010
Wd = freqNatural
v[0] = v0    #Sets the initial velocity
y[0] = y0    #Sets the initial position
```



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Simulation loop

```
for i in range(0, steps):  
    v[i+1] = (-k/m)*y[i]*dt + v[i]*(1-dt*c/m) + F0/m*np.cos(Wd*i*dt)*dt  
    y[i+1] = y[i] + v[i+1]*dt
```

```
fig, ax = plt.subplots()  
ax.plot(t, y, '-g', label='Undampened')
```

Damping set to 10% of critical damping

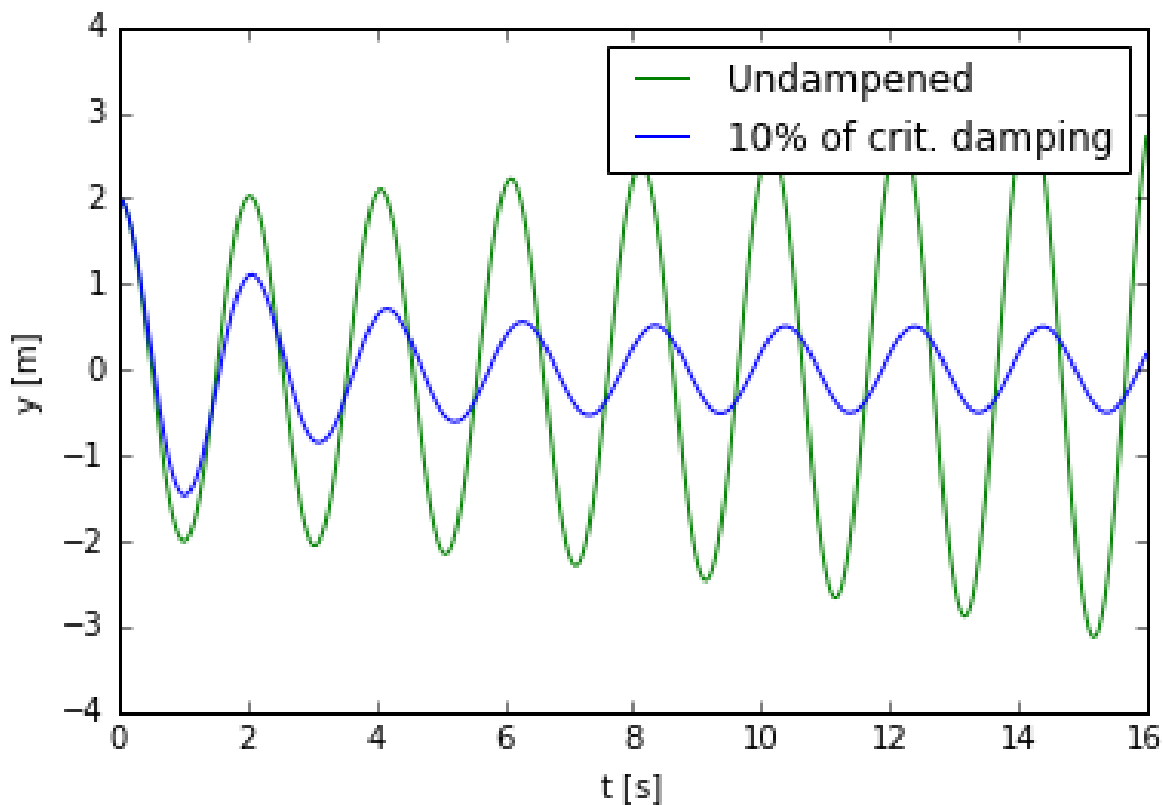
```
c = (2*(k*m)**0.5)*0.10
```

```
for i in range(0, steps):  
    v[i+1] = (-k/m)*y[i]*dt + v[i]*(1-dt*c/m) + F0/m*np.cos(Wd*i*dt)*dt  
    y[i+1] = y[i] + v[i+1]*dt
```

```
ax.plot(t, y, 'b-', label = '10% of crit. damping')  
ax.xlabel('t [s]')  
ax.ylabel('y [m]')  
ax.legend(loc = 'upper right')  
plt.show()
```

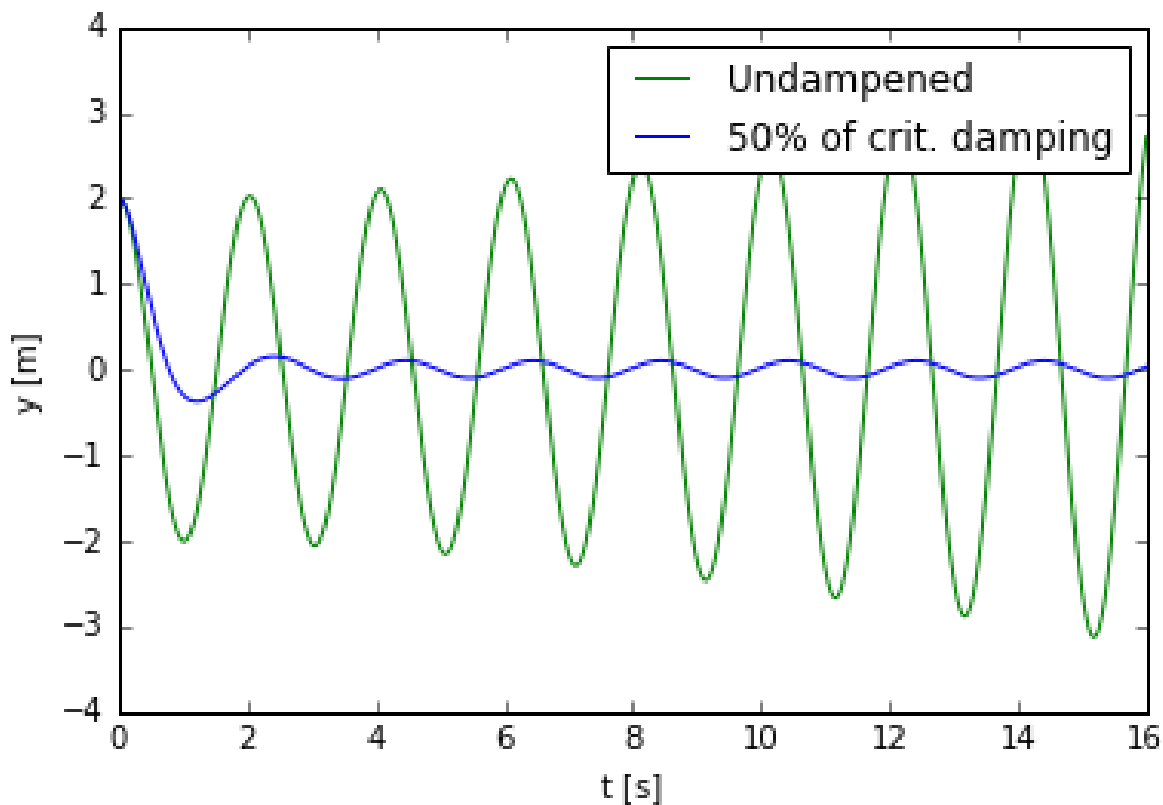


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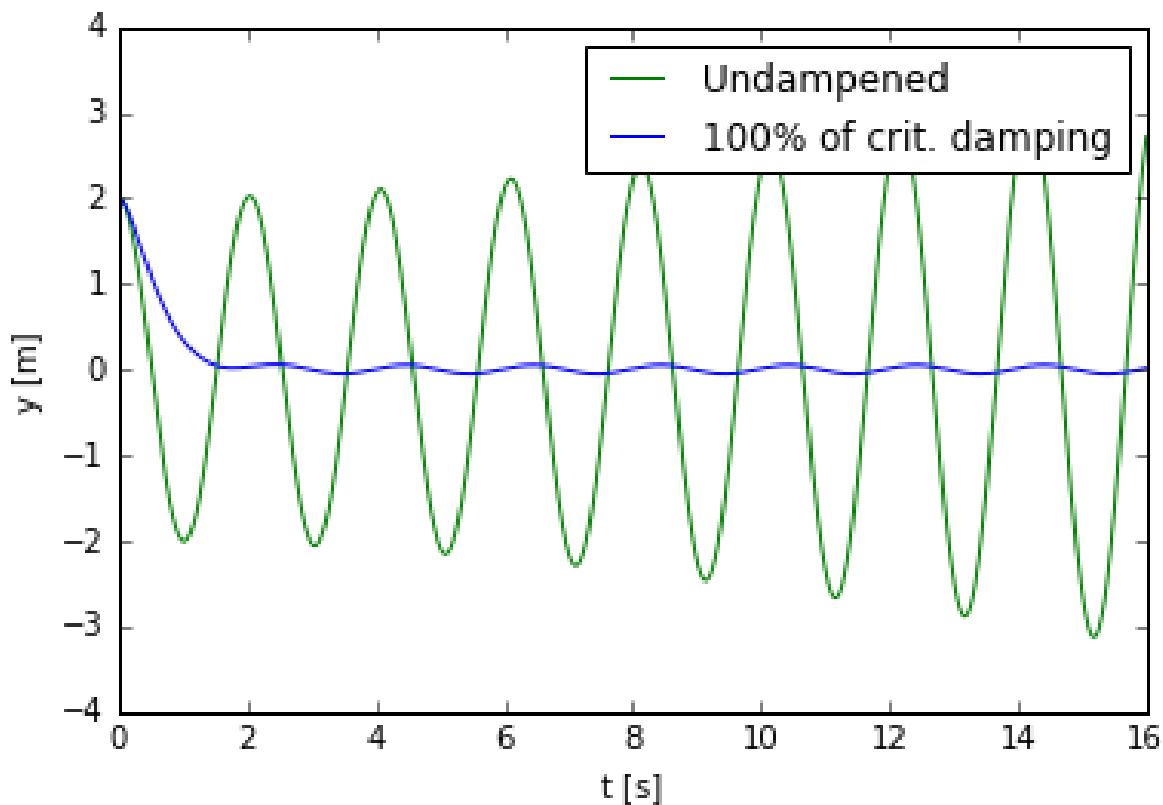


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Viscoelastic Materials

- Viscosity = The resistance to motion of a liquid
- Viscoelastic materials

Properties = Elastic solid + Fluid

= Materials that have mechanical properties dependent on time (loading rate or strain rate) and temperature



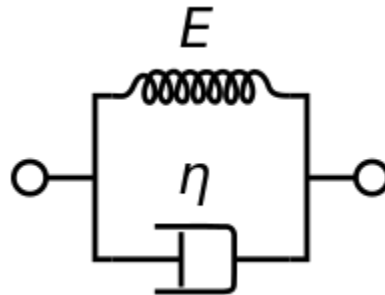


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Material Modelling

- Kelvin–Voigt material

A Kelvin–Voigt material, also called a Voigt material, is a viscoelastic material having the properties both of elasticity and viscosity. It is named after the British physicist and engineer Lord Kelvin and after German physicist Woldemar Voigt. the Kelvin–Voigt model does not describe stress relaxation.



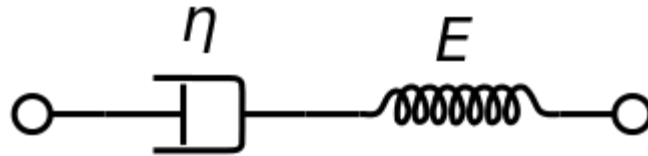


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Material Modelling

- Maxwell material

A Maxwell material is a viscoelastic material having the properties both of elasticity and viscosity. It is named for James Clerk Maxwell who proposed the model in 1867. It is also known as a Maxwell fluid. The Maxwell model does not describe creep or recovery.



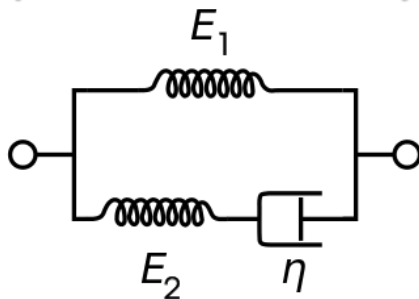


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Material Modelling

- Standard linear solid material

The standard linear solid (SLS) model is a method of modeling the behavior of a viscoelastic material using a linear combination of springs and dashpots to represent elastic and viscous components, respectively. SLS is the simplest model that predicts creep or recovery, and stress relaxation phenomena.





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scipy.integrate.solve_ivp

```
scipy.integrate.solve_ivp(fun, t_span, y0, method='RK45', t_eval=None, dense_output=False, events=None, vectorized=False, args=None, **options)
```

[\[source\]](#)

Solve an initial value problem for a system of ODEs.

This function numerically integrates a system of ordinary differential equations given an initial value:

Simulation loop

```
for i in range(n-1):
```

```
    F = -k*y[i] - m*g
```

```
    a = F/m
```

```
    v[i+1] = v[i] + a*dt
```

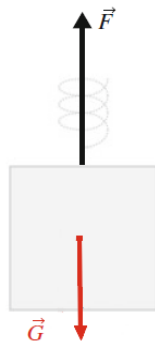
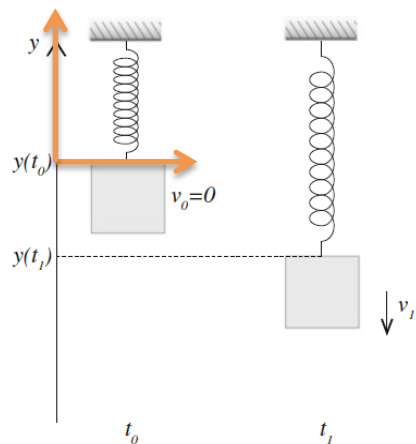
```
    y[i+1] = y[i] + v[i+1]*dt
```

```
    t[i+1] = t[i] + dt
```

Euler integration



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$$F_{net} = \vec{F} - \vec{G}$$

$$ma = -ky - mg$$

$$a = -ky/m - g$$

```
def MassSpring(t, state, m, k, g):  
    # unpack the state vector  
    x = state[0]  
    xd = state[1]  
    # compute acceleration xdd  
    F = -k*x - m*g  
    xdd = F/m    # xdd = ((k*x)/m) - g  
    #return the two state derivatives  
    return [xd, xdd]
```



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```
import numpy as np
import matplotlib.pyplot as plt
from scipy.integrate import solve_ivp

# simulation parameters
stime = 20.0 # time [s]
m = 1.0      # mass [kg]
k = 100.0    # spring stiffness [N/m]
v0 = 10.0    # initial velocity [m/s]
g = 9.8      # gravity [m/s^2]
dt = 0.5     # time step [s]
n = int(round(stime/dt))
tspan = np.linspace(0.0, stime, n)
```



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ODE function

```
def MassSpring(t, state, m, k, g):
```

```
    x = state[0]
```

```
    xd = state[1]
```

```
    # compute acceleration xdd
```

```
    F = -k*x - m*g
```

```
    xdd = F/m      # xdd = ((k*x)/m) - g
```

```
    # return the two state derivatives
```

```
    return [xd, xdd]
```

```
sol = solve_ivp(MassSpring, [tspan[0], tspan[-1]],
```

```
                0, v0],
```

```
                method = "RK45",
```

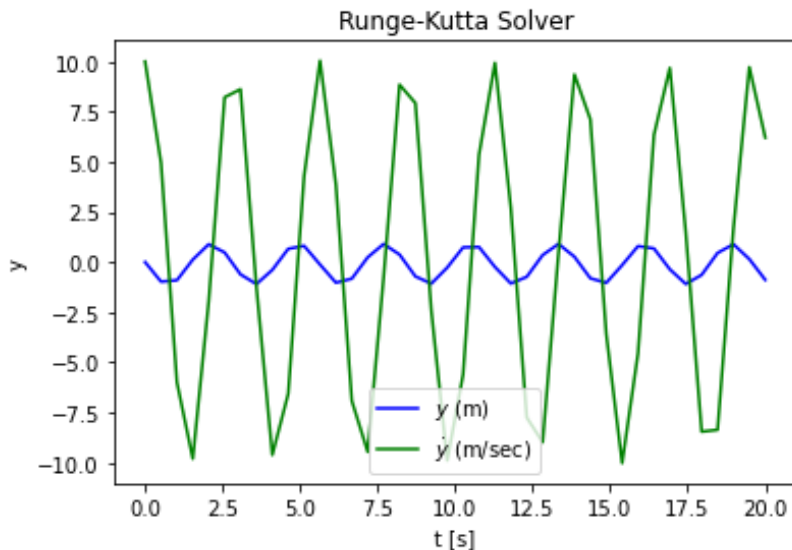
```
                t_eval=tspan, args=(m, k, g))
```



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```
# %% Plot states
```

```
fig, ax = plt.subplots()
ax.plot(sol.t, sol.y[0], '-b')
ax.plot(sol.t, sol.y[1], '-g')
ax.set_xlabel('t [s]')
ax.set_ylabel('y')
plt.title('Runge-Kutta Solver')
ax.legend(('y (m)', 'y (m/sec)'))
plt.show()
```





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'RK45' (default): Explicit Runge-Kutta method of order 5(4). The error is controlled assuming accuracy of the fourth-order method, but steps are taken using the fifth-order accurate formula.

'RK23': Explicit Runge-Kutta method of order 3(2). The error is controlled assuming accuracy of the second-order method, but steps are taken using the third-order accurate formula.

'DOP853': Explicit Runge-Kutta method of order 8. Python implementation of the "DOP853" algorithm originally written in Fortran.

'Radau': Implicit Runge-Kutta method of the Radau IIA family of order 5. The error is controlled with a third-order accurate embedded formula.

'BDF': Implicit multi-step variable-order (1 to 5) method based on a backward differentiation formula for the derivative approximation.

'LSODA': Adams/BDF method with automatic stiffness detection. This is a wrapper of the Fortran solver from ODEPACK.



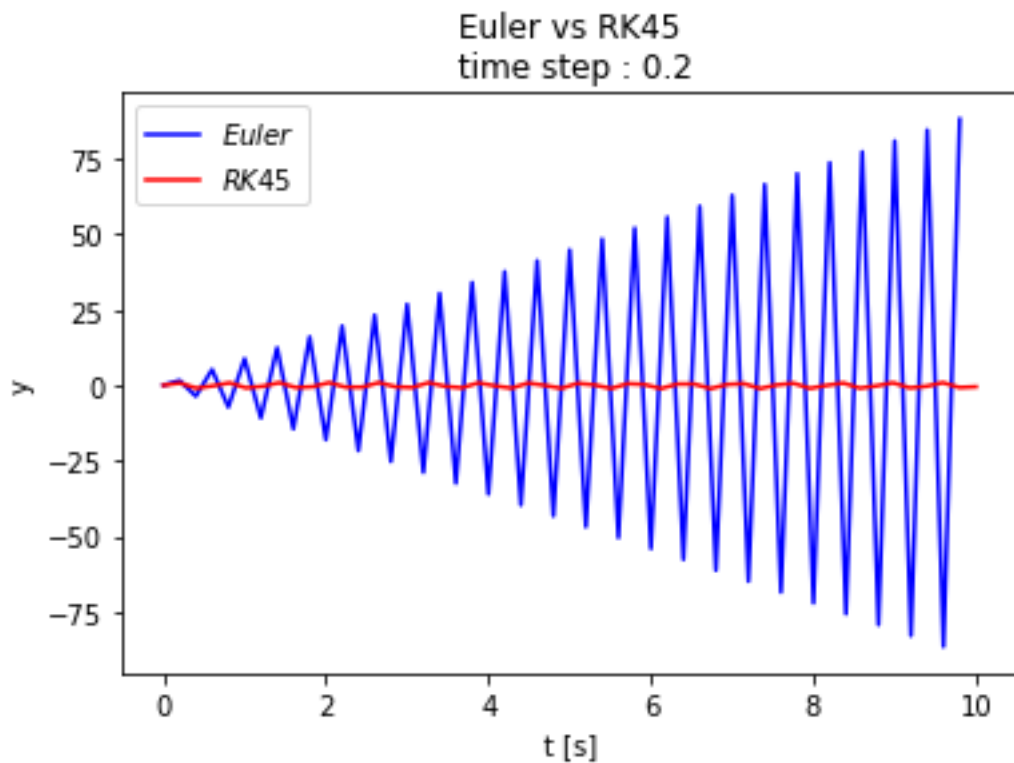
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Explicit Runge-Kutta methods ('RK23', 'RK45', 'DOP853') should be used for non-stiff problems and implicit methods ('Radau', 'BDF') for stiff problems. Among Runge-Kutta methods, 'DOP853' is recommended for solving with high precision (low values of rtol and atol).

If not sure, first try to run 'RK45'. If it makes unusually many iterations, diverges, or fails, your problem is likely to be stiff and you should use 'Radau' or 'BDF'. 'LSODA' can also be a good universal choice, but it might be somewhat less convenient to work with as it wraps old Fortran code.

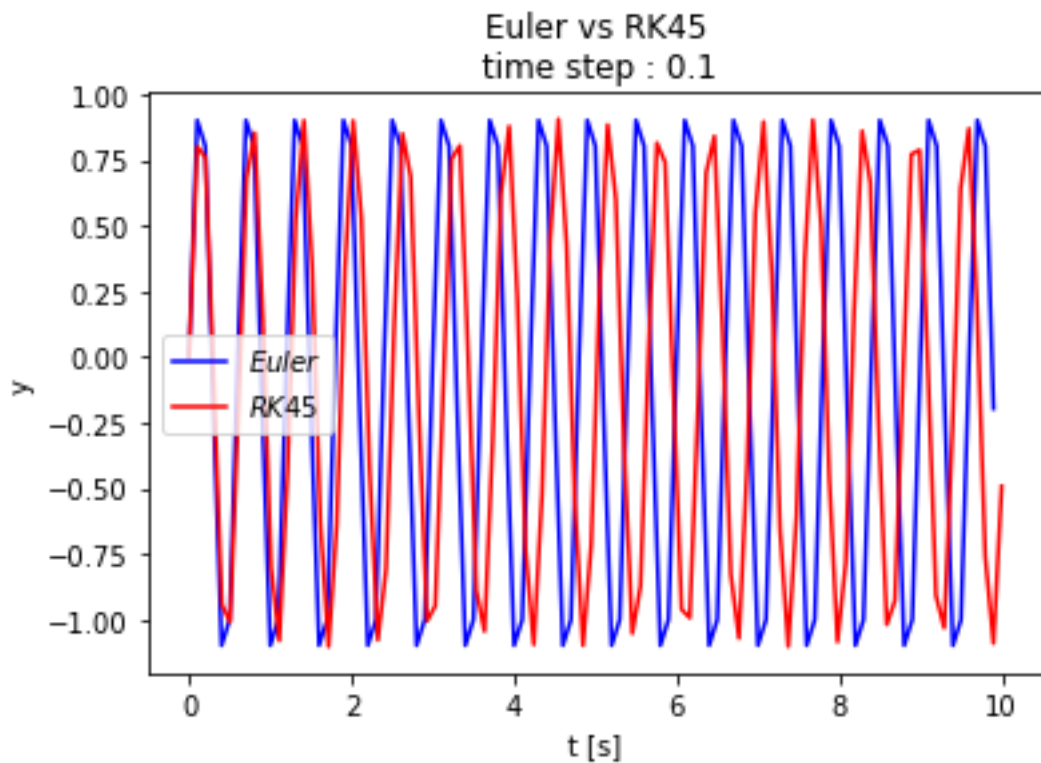


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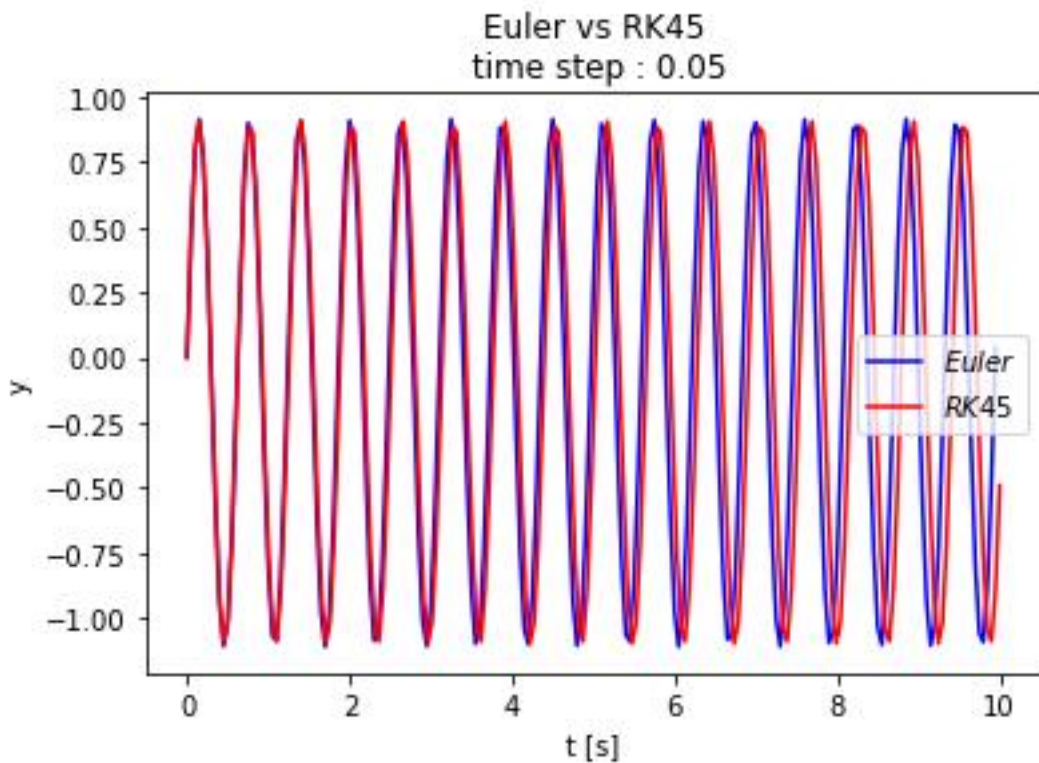


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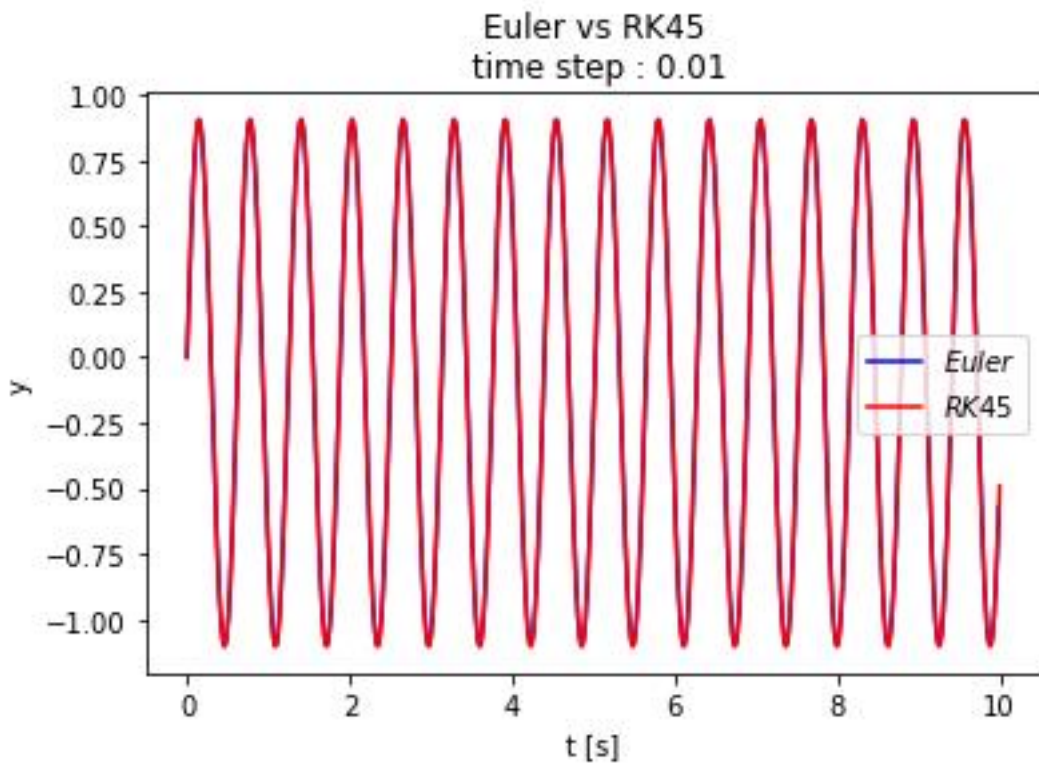


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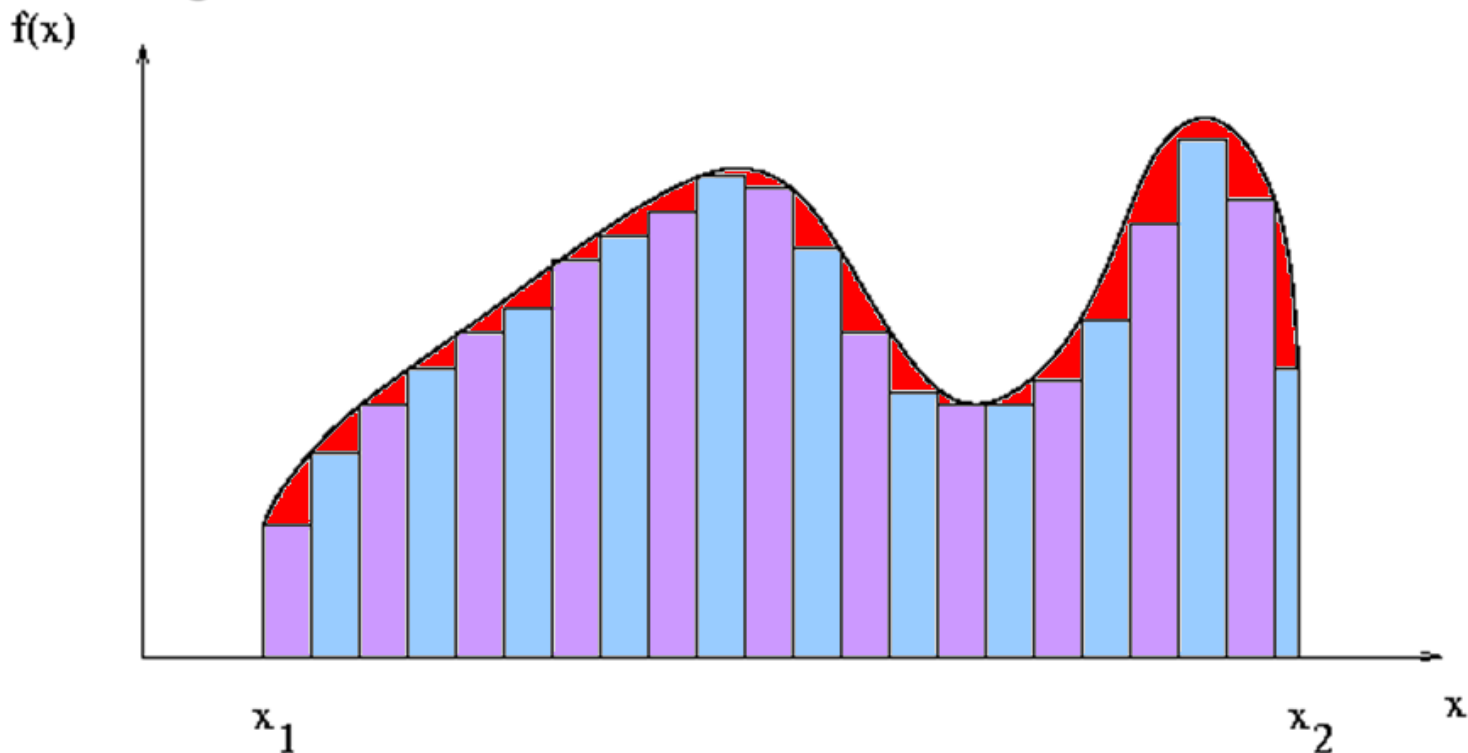
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Numerical Integration





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Euler : Euler's method is first order method. It is a straight-forward method that estimates the next point based on the rate of change at the current point. It is a single step method. If no dampening is used, particles get more and more energy over time. “**For example, bouncing particles will bounce higher and higher each time**”. Use this integrator for short simulations or simulations with a lot of dampening where speedy calculations are more important than accuracy.

Midpoint : Also known as “**2nd order Runge-Kutta**”. Slower than Euler but much more stable.

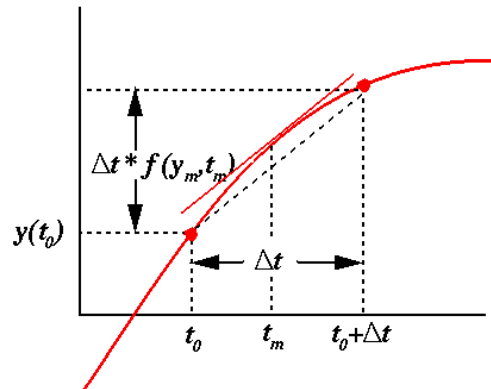
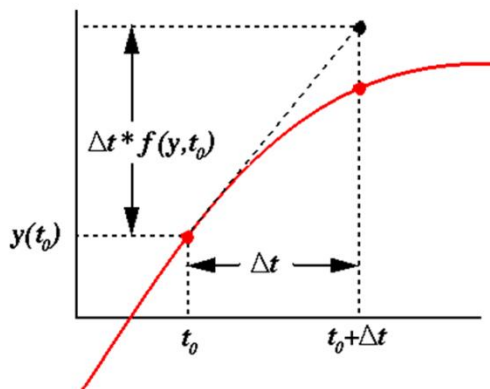
RK4 : Short for “**4th order Runge-Kutta**”. Similar to Midpoint but slower and in most cases more accurate. It is energy conservative even if the acceleration is not constant.



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Euler method

Also known as “**Forward Euler**”. Simplest integrator. Very fast but also with less exact results. If no damping is used, particles get more and more energy over time. Notably, Forward Euler's method is unconditionally unstable for un-damped oscillating systems (such as a spring-mass system or wave equations) in space discretization.





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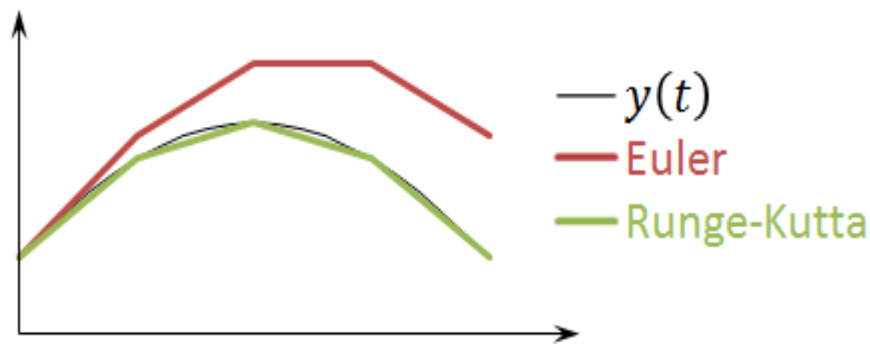
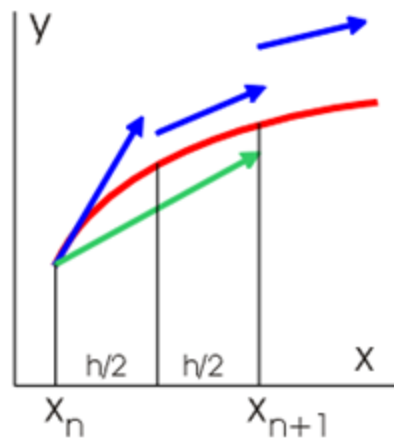
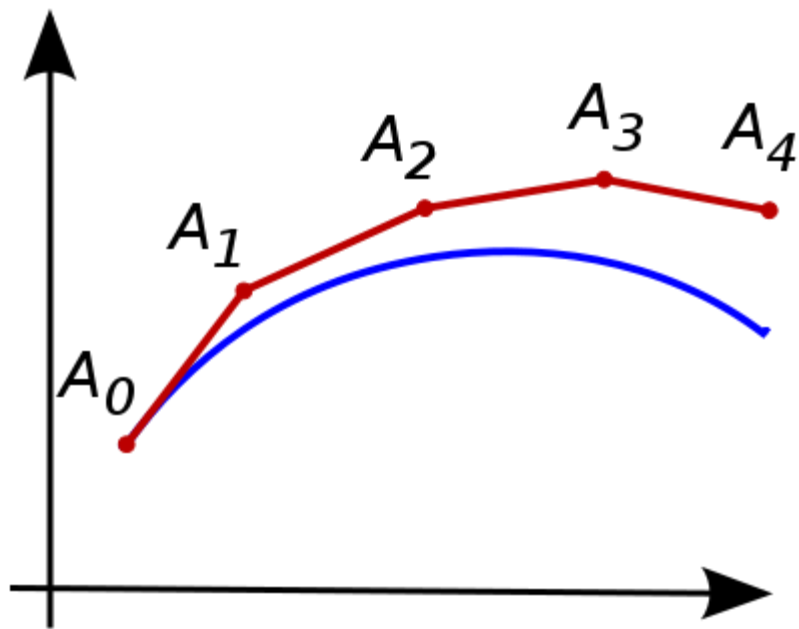
Euler method

Adding physical damping to the model (e.g. Rayleigh damping). In this case the damping will only be applied to the structure and the question is how much damping to introduce when physically it is negligible.

Adding numerical damping. This reduces the numerical oscillations, but also reduces the physical response which should be solved for, and the question is how much numerical damping to introduce in order to obtain acceptable results.



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Carl David Tolme **Runge** – Martin Wilhelm **Kutta**



*Carl David
Tolmé Runge*

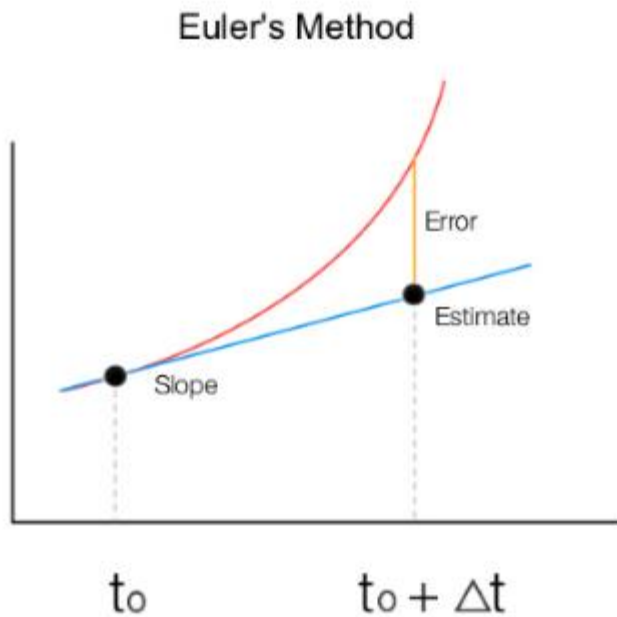


Martin Wilhelm Kutta



PHYSICS in COMPUTER ANIMATIONS and GAMES

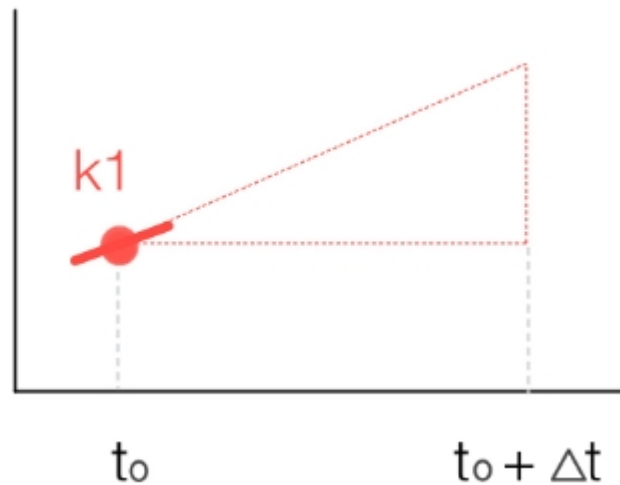
Runge –Kutta





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Runge –Kutta

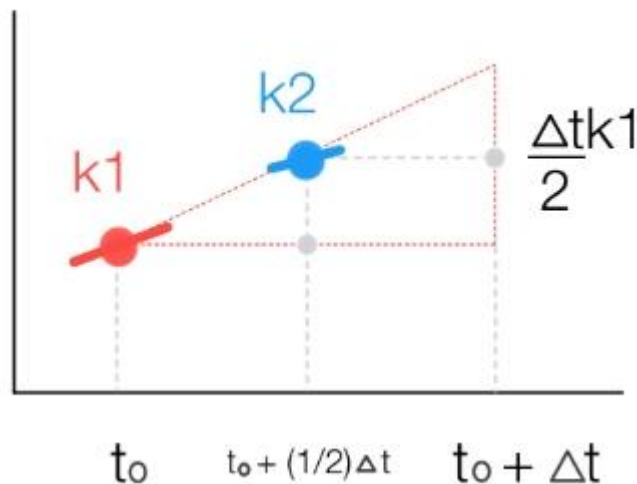


- 1) At time interval t_0 , calculate slope $k1$.
- 2) Create a triangle by projecting $k1$ to $t_0 + \Delta t$.



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Runge –Kutta

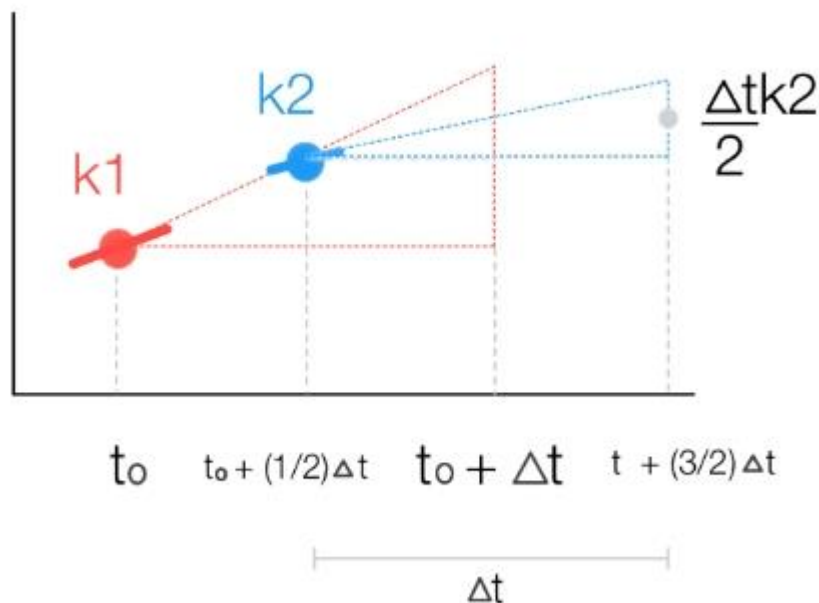


- 3) Calculate half the height of the triangle, $(\Delta t/2) k1$, and draw a horizontal line.
- 4) Draw a vertical line at interval $t_0 + (1/2)\Delta t$.
- 5) At this intersection point, calculate the slope $k2$.



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Runge –Kutta

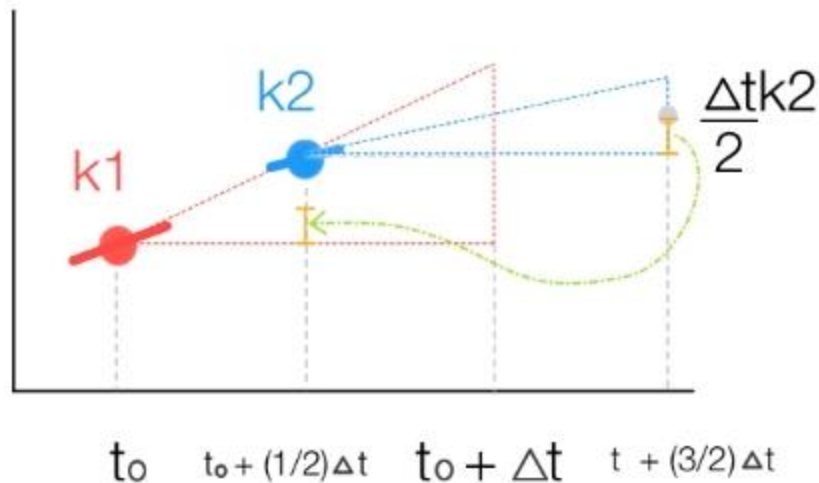


- 6) Starting out with $k2$, create a triangle by projecting $k2$ to $t + (3/2)\Delta t$.
- 7) Calculate half the height of the triangle, $(\Delta t/2)k2$.



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Runge –Kutta

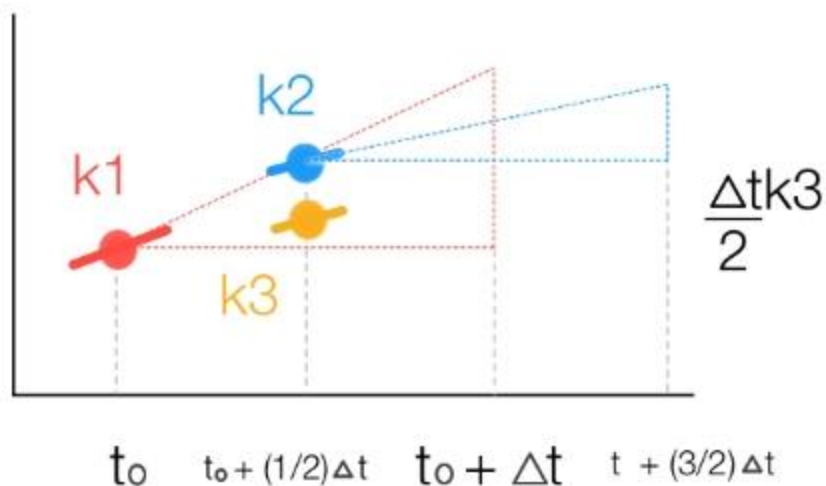


8) Translate the height to the base of the triangle formed by **k1** at $t_0 + (1/2)\Delta t$.



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Runge –Kutta

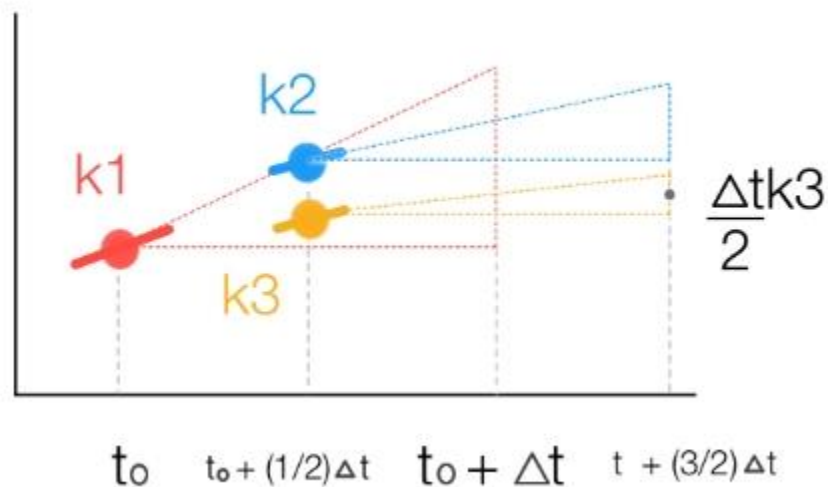


9) At this point, calculate the slope k_3 .



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Runge –Kutta



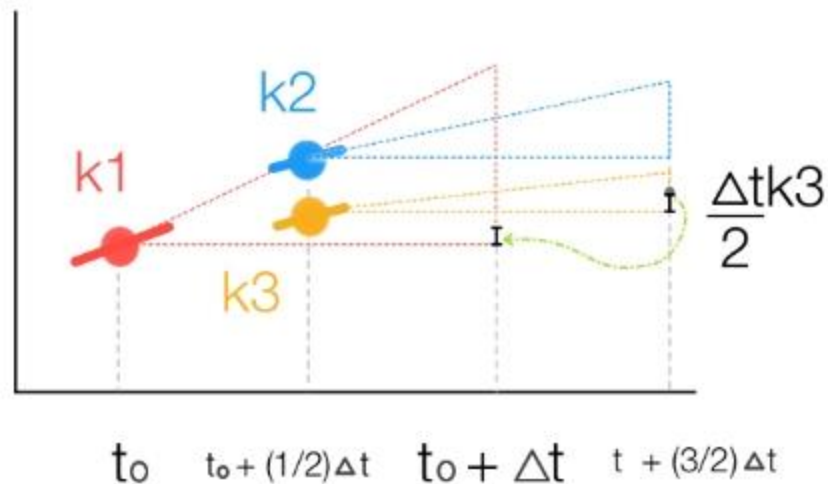
10) Create the another triangle by projecting k_3 to $t_0 + (3/2)\Delta t$

11) Find half the height of the triangle, $(\Delta t/2)k_3$



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Runge –Kutta

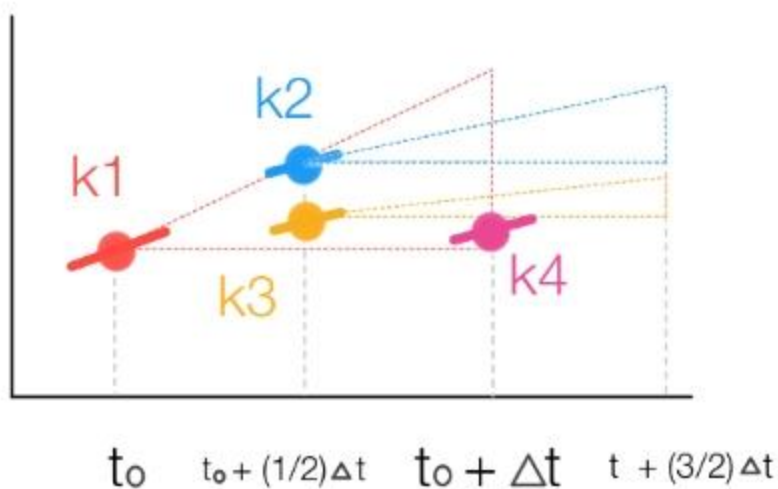


12) Translate this height to the base of the triangle formed by $k1$ at point $t_0 + \Delta t$



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Runge –Kutta



13) Find the slope k_4 .



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Runge –Kutta

The Runge-Kutta uses these slopes as weighted average to better approximate the actual slope, velocity, of the object.

Slopes

$$v_{t+\Delta t} = v_t + \frac{\Delta t}{6} (k_1 + 2k_2 + 2k_3 + k_4)$$

The diagram illustrates the Runge-Kutta method's use of four slopes (k1, k2, k3, k4) to approximate the next velocity value. The word "Slopes" is written above the equation, with four arrows pointing down to each of the slope terms: k1 (red), k2 (blue), k3 (yellow), and k4 (pink). The equation shows these slopes being weighted (k2 and k3 are multiplied by 2) and summed, then multiplied by the time step Δt/6, and added to the current velocity v_t to find the next velocity v_{t+Δt}.



PHYSICS in COMPUTER ANIMATIONS and GAMES

How to set an icon in Pygame.

```
import pygame
```

```
WHITE = (255, 255, 255)
pygame.init()
icon = pygame.image.load('icon.png')
pygame.display.set_icon(icon)
size = (700, 200)
screen = pygame.display.set_mode(size)
pygame.display.set_caption("Physics in Game and Animation")
# Loop until the user clicks the close button.
done = False
# Used to manage how fast the screen updates
clock = pygame.time.Clock()
# ----- Main Program Loop -----
while not done:
    # --- Main event loop
    for event in pygame.event.get():
        if event.type == pygame.QUIT:
            done = True

    screen.fill(WHITE)
    pygame.display.flip()
    clock.tick(60)
# Close the window and quit.
pygame.quit()
```



Physics in Game and Animation

